

High-oxygen-pressure Crystal Growth of Ferroelectric $\text{Bi}_4\text{Ti}_3\text{O}_{12}$ Single Crystals

Ferroelectric bismuth titanate ($\text{Bi}_4\text{Ti}_3\text{O}_{12}$, BiT) has been regarded as a promising material for innovative semiconductor-based applications such as nonvolatile memories, electro-optic devices and uncooled infrared detectors because of its high Curie temperature, large spontaneous polarization (P_s) and large electro-optic coefficient [1]. The control of polarization states is the underlying basis of these functional devices, and polarization switching is achieved through the nucleation of domains and the following domain-wall motion by applying an electric field (E). Leakage current arising from defects, however, interferes with the polarization switching of BiT-based materials [2]. In addition, oxygen vacancies are known to act as an obstacle to the polarization switching, and a resultant remanent polarization (P_r) is suppressed by the clamping of the domain walls by oxygen vacancies [2]. The leakage current and domain clamping by oxygen vacancies make BiT unsuitable for the practical applications. Thus, a guiding principle of defect control is required to be established for obtaining high-quality BiT-based devices with a large P_r as well as a low leakage current. Here, we show that high-oxygen-pressure crystal growth is an effective process for obtaining high-quality BiT crystals with a large P_r and a low coercive field (E_c) as well as low leakage current [3].

Synchrotron radiation powder diffraction experiments on the crushed powder of the crystals were performed using a large Debye-Scherrer camera installed at **BL02B2** to investigate the precise crystal structure of BiT synthesized by the processing method. We used high-energy SR with a wavelength of $\lambda = 0.035639(2)$ nm ($E \sim 35$ keV) to reduce absorption by the samples. The BiT crystals under different P_{O_2} atmospheres had almost the same lattice

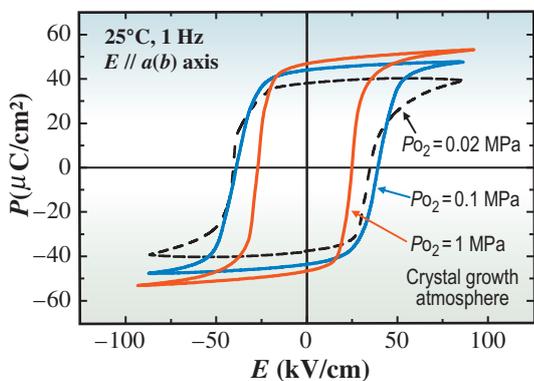


Fig. 1. Polarization hysteresis loops along the $a(b)$ -axis for the BiT crystals grown at a P_{O_2} of 0.02, 0.1, and 1 MPa. These crystals were annealed at 900°C for 10 h in air. The measurements were conducted at 25°C using an E at a frequency of 1 Hz.

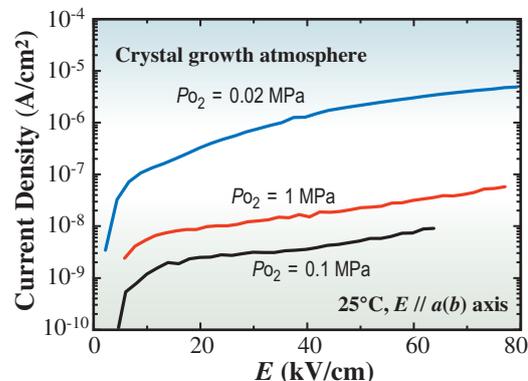


Fig. 2. Remanent polarization (P_r) and coercive field (E_c) as a function of P_{O_2} during crystal growth (25°C , 1 Hz). The values of P_r and E_c are average in positive and negative regions in the polarization hysteresis loops (Fig. 1).

parameters: $a = 0.54505(5)$ nm, $b = 0.54108(4)$ nm, $c = 3.2834(3)$ nm. Rietveld analyses demonstrated that there was no significant difference in crystal structure in crystals grown under different pressures [3]. These experimental results indicate that high-oxygen-pressure crystal growth is suitable for obtaining defect-controlled ferroelectric crystals without any significant change in the main crystal structure.

Domain observations by piezoresponse force microscopy (PFM) demonstrate that the clamping of 90° domains deteriorates P_r for the crystals grown at 0.02 MPa oxygen, which is suggested to originate from the strong attractive interaction between 90° domain walls and oxygen vacancies. The vacancy formation of Bi and O during crystal growth at high temperatures is suppressed at a higher oxygen pressure, leading to a larger P_r of $47 \mu\text{C}/\text{cm}^2$ and a lower E_c of 26 kV/cm for the crystals grown at 1 MPa oxygen [3].

Figure 1 shows the polarization hysteresis loops measured along the $a(b)$ -axis (25°C , 1 Hz) [3]. The crystals ($P_{\text{O}_2} = 0.02$ MPa) exhibited hysteresis with $P_r = 38 \mu\text{C}/\text{cm}^2$ and an $E_c = 38$ kV/cm. The high- P_{O_2} -grown crystals had larger values of P_r of $44 \mu\text{C}/\text{cm}^2$ ($P_{\text{O}_2} = 0.1$ MPa) and $47 \mu\text{C}/\text{cm}^2$ ($P_{\text{O}_2} = 1$ MPa). Note that the crystals grown at $P_{\text{O}_2} = 1$ MPa exhibited well-saturated polarization hysteresis with $E_c = 26$ kV/cm. This E_c value was much lower than those of the other crystals. Figure 2 shows P_r and E_c as a function of P_{O_2} during crystal growth [3]. With increasing P_{O_2} , P_r monotonically increased, while the decrease in E_c was marked over $P_{\text{O}_2} = 0.1$ MPa.

Figure 3 shows the leakage current density (J) as a function of E along the $a(b)$ -axis (25°C) [3]. The

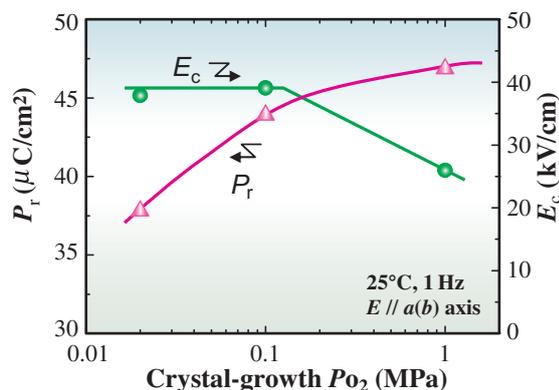


Fig. 3. Leakage current density (J) as a function of E along the $a(b)$ -axis for the BiT crystals (25°C, 1 Hz).

crystals ($Po_2 = 0.02$ MPa) exhibit a high J of the order of 10^{-7} to 10^{-6} A/cm². The increase in Po_2 to 0.1 MPa led to a drastic decrease in J to the order of 10^{-9} A/cm². The crystals grown at $Po_2 = 1$ MPa exhibited a relatively low J of the order of 10^{-8} A/cm².

Here, we discuss the mechanism of domain clamping during the polarization switching along the $a(b)$ -axis. Figure 4 shows the PFM images observed on the $a(b)$ - c surface of the crystals grown at $Po_2 = 0.02$ MPa [3]. The in-plane PFM image of the as-annealed (nonpoled) crystals (Fig. 4(a)) exhibits 180° DWs parallel to the a - b plane. After an E of 100 kV/cm was applied along the $a(b)$ -axis at 25°C, the out-of-plane PFM image of the poled crystals was observed (Fig. 4(b)). A single domain state was not established for the poled crystals even though the applied E (100 kV/cm) is much higher than the E_c value (38 kV/cm). Domains with $P_{s(a)}$ parallel to the poling direction were found. This is direct evidence that 90° domains are switched by applying an E of 100 kV/cm. Note that unswitched regions, i.e., 90° domains with $P_{s(a)}$ normal to the poling direction remained, and 90° DWs with an irregular structure

appeared. The irregular-shaped 90° DWs have been reported to originate from the attractive interaction between $V_O^{..}$ and the electric field established near the 90° DWs due to the discontinuity of the P_s component normal to the DWs. In the domains with $P_{s(a)}$ parallel to the poling direction, a small number of 180° domains with $P_{s(a)}$ antiparallel to the poling direction were observed. These 180° domains are a result of the domain backswitching due to the depolarization field. Our PFM observations lead to the conclusion that the clamping of 90° DWs plays a detrimental role in the $P_{s(a)}$ polarization switching in the BiT crystals. The vacancy formation at high temperatures is suppressed under a higher- Po_2 atmosphere, and then $[V_O^{..}]$ becomes lower for the crystals grown at a higher Po_2 . The larger P_r observed for the crystals ($Po_2 = 1$ MPa) is found to originate from suppressed 90° domain clamping because of a lower $[V_O^{..}]$.

In summary, the effects of Po_2 during the crystal growth of BiT on domain-switching behavior have been investigated through polarization measurements and domain observations by PFM. The crystal structure is investigated by high energy synchrotron radiation powder diffraction. The crystals grown at a high Po_2 of 1 MPa showed a large P_r of 47 $\mu\text{C}/\text{cm}^2$ and a low E_c of 26 kV/cm. PFM observations demonstrate that the clamping of 90° DWs plays a detrimental role in polarization switching, leading to a low P_r . High- Po_2 sintering is proposed as an effective process for suppressing the formation of vacancies of Bi and O without any change in the main crystal structure, leading to the realization of high-quality BiT-based devices with enhanced polarization-switching properties as well as low leakage current.

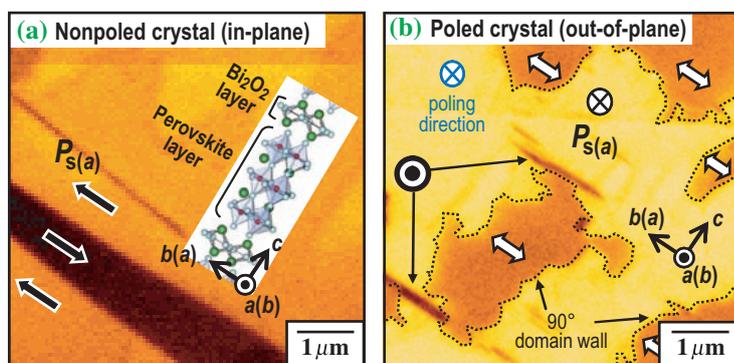


Fig. 4. PFM images of the $a(b)$ - c surface of the BiT crystals grown in air ($Po_2 = 0.02$ MPa); (a) in-plane image of the nonpoled (as-annealed) crystal, and (b) out-of-plane image of the poled crystal. The poling was conducted by applying an E of 100 kV/cm along the $a(b)$ -axis at 25°C.

Yuji Noguchi^a, Masaru Miyayama^a
and Yoshihiro Kuroiwa^b

^a Research Center for Advanced Science and Technology, The University of Tokyo

^b Department of Physical Science, Hiroshima University

*E-mail: ynoguchi@crm.rcast.u-tokyo.ac.jp

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