

Symmetry of wavefunction in perpendicular magnetic anisotropy films

Recent high-density magnetic storage has been made possible by perpendicular magnetic recording (PMR) techniques. The first application of PMR to hard disk drives was achieved in Dec. 2005 [1]. PMR can achieve 5-10 times higher recording density than the previous conventional magnetic recording. The key material for PMR is a film with perpendicular magnetic anisotropy (PMA), which has a large positive magnetic anisotropy energy perpendicular to the film plane.

Metallic multilayers, such as Co/Pd and Co/Pt, exhibit large PMA energies when their magnetic layer thickness is reduced to a few monolayers [2-5]. In spite of many experimental and theoretical studies toward the understanding of magnetic anisotropy, the origin of PMA in multilayered magnetic thin films remains to be clarified.

Previous theoretical investigations [6,7] have shown the following. The orbital hybridization between Co-3d and Pd-4d (Pt-5d) at the interface increases the population of symmetric Co-3d states with $|m|=2$. Here, m denotes the magnetic quantum number. The angular momentum induced by the increased number of $|m|=2$ states gives enhanced orbital magnetic moments perpendicular to the film plane. The enhanced orbital magnetic moments perpendicular to the film plane induce PMA. X-ray circular magnetic dichroism (XMCD) experiments have shown the presence of enhanced Co 3d orbital magnetic moments perpendicular to the film plane at the Co/Pd (or Pt) interface [8,9]. However, the contributions of the $|m|=2$ symmetry to PMA have not yet been confirmed experimentally.

A magnetic Compton profile (MCP), $J_{mag}(p_z)$, is expressed by a projection of a spin density map to the p_z axis in momentum space. Here, p_z denotes the z component of the electron momentum \mathbf{p} in a solid. Because the symmetries of a wavefunction are the same in real space and momentum space, the symmetry of the spin-dependent wavefunction can be determined if MCPs are observed from different directions [10].

Recently, we have applied MCPs to the analysis of Co/Pd and Co/Pd multilayers to directly observe the $|m|=2$ symmetry of Co 3d states in multilayers with PMA. In this paper, we report the symmetry of the wavefunction in multilayers with PMA [11].

Five multilayer films (Co(0.8 nm)/Pd(x nm), $x = 0.8, 1.6, 4.0$, Co(0.8 nm)/Pt(x nm) $x = 0.8, 4.0$) were fabricated on PET film substrates of 4 μm thickness by RF sputtering. The total film thickness of the

multilayers was adjusted to about 1 μm . The thin-film samples were folded 16 times to increase their effective thickness, and the effective thicknesses of the films and PET substrate were 16 μm and 64 μm , respectively.

The crystal structure was confirmed by $\theta-2\theta$ X-ray diffraction measurement. The (111) texture of the fcc structure was observed in the intermediate-angle region. Satellite peaks, which were observed around the intermediate-angle region, confirm the designed period of the multilayers. The lattice constant obtained from the peak position of (111) has a linear relation with the thickness fraction of the Co layer.

Magnetization was measured in an out-of-plane configuration (applied fields were perpendicular to the sample surface) and an in-plane configuration (applied fields were parallel to the sample surface). Figure 1 shows the lattice constants and PMA energies. The present results reproduce those of previous studies [2-5]. Magnetic anisotropy changes from in plane magnetic anisotropy to PMA at a lattice constant of more than 0.372 nm.

MCPs were measured at the high-energy beamline BL08W. The circularly polarized X-ray energy was set to be 174 keV. The degree of circular polarization was about 0.76. The scattered X-rays were detected by a 10-segmented Ge solid-state detector (SSD) with a scattering angle of 178° that was installed 1 m from the sample. The momentum resolution was 0.43 atomic units (a.u.). The applied magnetic field, which was supplied by a superconducting magnet, was ± 2.5 T for magnetization saturation in both the in-plane and out-of-plane configurations. All the measurements were carried out under vacuum at room temperature.

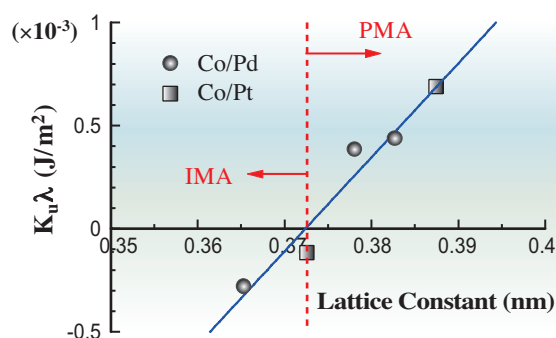


Fig. 1. Perpendicular magnetic anisotropy (PMA) energy density per bilayer against the lattice constants of fcc Co/Pd and Co/Pt multilayers. The blue solid line denotes a least squares fitting result as a visual guide. The red vertical dashed line denotes the boundary of PMA and in-plane magnetic anisotropy (IMA).

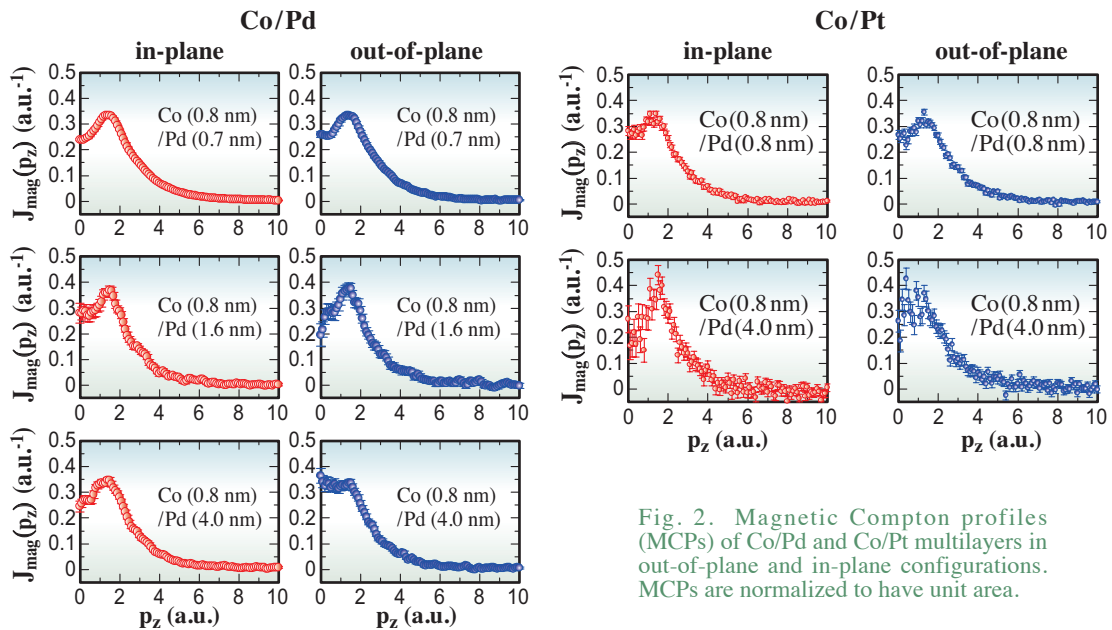


Fig. 2. Magnetic Compton profiles (MCPs) of Co/Pd and Co/Pt multilayers in out-of-plane and in-plane configurations. MCPs are normalized to have unit area.

Figure 2 shows the MCPs of the (Co(0.8 nm)/Pd(x nm) and (Co(0.8 nm)/Pt(x nm) multilayers with the in-plane and out-of-plane configurations. Differences in the MCP between the in-plane and out-of-plane configurations, which we call the anisotropies of the MCP, are observed for the Co/Pd and Co/Pt multilayers within the momentum region 2 a.u. The anisotropies depend on the Pd and Pt thickness, x . This dependence originates from the changes in the populations of $l_{m_l} = 0, 1$, and 2 states, and hence from the change in the symmetry of the wavefunctions. Then the MCPs are decomposed using the model calculation, and the populations of $l_{m_l} = 0, 1$, and 2 states are obtained.

Figure 3 shows the lattice constants and populations of $l_{m_l} = 0, 1$, and 2 states. $l_{m_l} = 2$ states

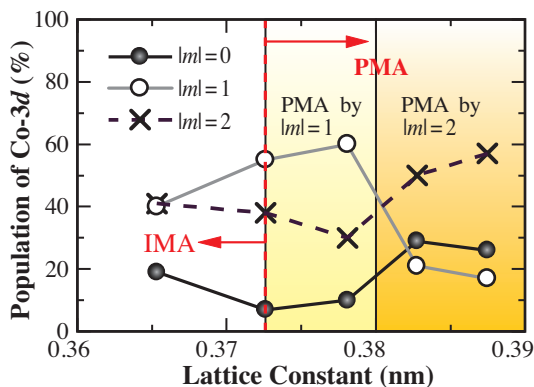


Fig. 3. Population of $l_{m_l} = 0, 1$, and 2 states in the spin-projected Co 3d electrons against the lattice constants of fcc Co/Pd and Co/Pt multilayers. The red vertical dashed line denotes boundary of PMA and in-plane magnetic anisotropy (IMA).

dominate above 0.380 nm. This indicates that $l_{m_l} = 2$ states contribute to PMA. $l_{m_l} = 1$ states dominate between 0.372 nm and 0.380 nm. This indicates that $l_{m_l} = 1$ states also contribute to PMA on the boundary region.

In conclusion, the lattice constants change the PMA energy and the symmetry of the wavefunction. The contributions of the $l_{m_l} = 2$ symmetry to PMA are confirmed experimentally. The contributions of the $l_{m_l} = 1$ symmetry to PMA are also experimentally observed for the first time.

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