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## Pressure effect in metal-insulator transition system (Y,Ca)TiO<sub>3</sub> and Kondo insulator, CeRhAs and YbB<sub>13</sub>

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The rare earth compounds CeRhAs and YbB<sub>12</sub> which form an energy gap and show semiconducting behavior by K ondo effect, and the transition metal alloys (Y,Ca)TiO, which causes metal-insulator transition due to strong Coulomb correlation have attracted the attention of many researchers.

It is known that CeRhAs successively forms a super-lattice structure at T = 370, 235 and 165 K. respectively. The mechanism of an energy gap formation may have an relation to the lattice instability due to this super-lattice formation. In order to acquire the information about energy gap formation from the response of the lattice against pressure, we investigated the relation between a lattice constant and pressure by using a DAC, up to 2.5 GPa, and from T = 100 to 260 K. The temperature dependence of the lattice constant, a, of orthorhombic structure at various pressures is shown in Fig. 1. From this figure, at the temperature below  $T_3$ , we have found that the lattice constant changes from decreasing behavior to increasing with increasing pressure.

Fig. 2 shows pressure dependence of the lattice constant at T = 230 K of cubic Kondo semiconductor YbB<sub>12</sub>. The loaded pressure was up to 50 GPa. Observed hump around 30 GPa at 300 K by T.L. should be related f-state of YbB<sub>12</sub>, but strange hysteresis at this work is shown at 230 K.

Furthermore we have studied the pressure dependence of lattice constant of metal-

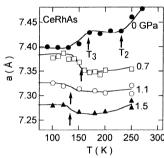


Fig. 1 Temperature dependence of lattice constant of CeRhAs at each pressure.

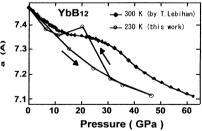


Fig. 2 Pressure dependence of latice constant of YbB<sub>12</sub> at T = 230 K.

insulator transition system Y<sub>0.63</sub>Ca<sub>0.37</sub>TiO<sub>3</sub> up to 46 GPa. Phase seperation was entirely disappeared at P > 2 GPa. The lattice volume became small with increase in pressure, and especially abnormalities were not seen in this variation.

## The process of iron nitrides syntheses under high pressure and temperature.

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Iron and nitrogen forms many kind of iron nitrides, in which interstitial nitrogen atoms influence upon the magnetic properties. Although iron nitrides of Fe<sub>4</sub>N, Fe<sub>3</sub>N, and Fe<sub>2</sub>N, can be synthesized from iron (bcc phase) and gaseous nitrogen at atmospheric pressure, there are composition limits in nitrogen content due to the difficulty in the interstitial of nitrogen atoms into bcc iron structure. We try to synthesize iron nitrides from high pressure phases of iron (hcp or fcc phase) and a supercritical fluid of nitrogen using a in-situ laser heating method under high pressure up to 30 GPa. In the previous study, we confirmed two kinds of iron nitride compounds for different pressure ranges by using a high pressure x-ray diffraction method and Moessbauer spectroscopy for laser-heated samples. One is known as a Fe<sub>3</sub>N compound, which appears below 10 GPa after a laser heating. The other is a new magnetic phase, which can be synthesized at high pressure above 15 GPa. However, it partially decompose to fcc phase at atmospheric pressure. The high pressure phase of hcp iron is stable above 15 GPa even at room temperature. On the other side, hcp iron

converts into fcc structure at high temperature. In the present study, we tried to check which phase reacts with nitrogen to produce the new magnetic phase.

An iron foil sample was placed into a rhenium gasket with liquid nitrogen in a DAC. After applying a pressure up to 30 GPa, a continuous Nd:YAG laser was focused on to the sample from both sides. The temperature was monitored by radiation spectra. During heating, the sample was exposed by a monochromatic x-ray at BL-10XU. The diffraction patterns were continuously taken by an imaging plate under high P, T condition. The hcp iron quickly changes into the new magnetic phase at around 1300 K, which coincides with a transition temperature to fcc iron. Furthermore, we could not obtain the diffraction pattern from the fcc iron. Consequently, we conclude that the formation of the new phase occurred on the phase boundary between fcc and hcp, and the fcc is a key phase to produce the new phase.